

of the reactions of **2** with phosphines indicate that a complicated mechanism, probably involving two phosphine-trappable intermediates in thermal equilibrium with **2**, is operative. Significantly, the reaction of **2** with phosphines is the first dimetallacycle transformation we have found in which the organometallic fragment is extruded in dinuclear form. This result provides a unique opportunity to employ crossover experiments to determine whether the two metal centers remain associated with each other during the reaction; i.e., whether the *o*-xylylene extrusion is truly a dinuclear elimination reaction. The outcome and interpretation of these experiments will be reported in a full paper.

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ESR Detection of the Dimethyl Ether Radical Cation

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Despite their great importance as reaction intermediates, comparatively few oxygen-centered radicals have been characterized by ESR spectroscopy.^{1,2} Here we report the radical cation of dimethyl ether, $(\text{CH}_3)_2\text{O}^+$, which is isoelectronic with the dimethyl aminyl radical $(\text{CH}_3)_2\text{N}^{\cdot}$.³⁻⁵ As far as we are aware, this is the first example of an oxygen-centered radical cation where the unpaired electron is not delocalized into an aromatic ring system as in the case of the anisole⁶ and dimethoxybenzene⁷ radical cations.

Bearing in mind the high reactivity of alkoxy radicals in hydrogen atom abstraction reactions,⁸ the dimethyl ether radical cation would be expected to react avidly with any potential hydrogen donor. This expectation is borne out by the very large rate constant of $1.9 \times 10^{-9} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ ($1.1 \times 10^{12} \text{ L mol}^{-1} \text{ s}^{-1}$) measured for the gas-phase ion-molecule reaction of $(\text{CH}_3)_2\text{O}^+$ with $(\text{CH}_3)_2\text{O}$ to form the dimethyloxonium cation $(\text{CH}_3)_2\text{OH}^+$,⁹ the reaction occurring essentially on every collision. Thus, $(\text{CH}_3)_2\text{O}^+$ would probably be difficult to detect if generated in a hydrogen-containing solvent or matrix. This problem has been avoided by using a γ -irradiation technique¹⁰⁻¹² which allows the radical cation to be generated through positive charge transfer from a Freon matrix to the dimethyl ether solute, the method satisfying the requirement that the $(\text{CH}_3)_2\text{O}^+$ species be trapped in a chemically inert environment.

Figure 1 shows the first-derivative and second-derivative ESR spectra obtained from a γ -irradiated solid solution of dimethyl

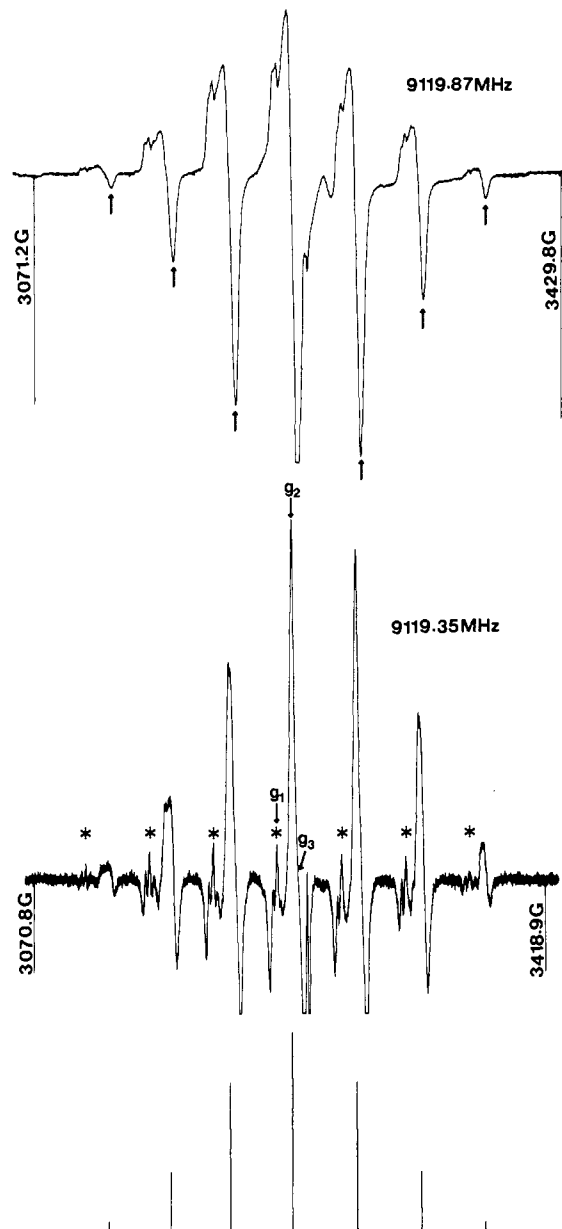


Figure 1. First-derivative (upper) and second-derivative (lower) ESR spectra of a γ -irradiated solid solution of 5 mol % dimethyl ether in trichlorofluoromethane recorded at 97 K after irradiation (dose, 1 M rad) at 77 K. The stick diagram shows the seven hyperfine components from the dimethyl ether radical cation, the anisotropy in the line positions extending between the two sets of features marked by asterisks (lower spectrum) and arrows (upper spectrum). The larger signal amplitudes for the high-field lines are due to a line-narrowing effect resulting from $g_1 > g_3$ and $A_1 > A_3$ (see text).

Table I. ESR Parameters for the Dimethyl Ether Radical Cation in a Freon 11 Matrix at 97 K

components of g tensor and g_{iso}	^1H hyperfine couplings, G
$g_1 = 2.0138$	$A_1 = 43.6$
$g_2 = 2.0072$	$A_2 = 42.8$
$g_3 = 2.0045$	$A_3 = 42.5$
$g_{\text{iso}} = 2.0085$	$a_{\text{B}}(6\text{H})^a = 43.0$

^a Isotropic coupling = $(1/3)(A_1 + A_2 + A_3)$.

ether in trichlorofluoromethane (Freon 11) at 97 K. The first-derivative spectrum consists of seven lines with approximately binomial intensity ratios (1:6:15:20:15:6:1) as expected for hyperfine interaction with six equivalent ^1H ($I = 1/2$) nuclei. Each line in the binomial pattern is asymmetrically broadened, and this

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Table II. Spin Densities and ^1H Hyperfine Couplings Calculated for the Dimethyl Ether Radical Cation by the INDO Method^a

atom	coordinates, ^a Å			atomic orbital	spin density	^1H hyperfine coupling, ^b G
	x	y	z			
C	0.983	0	0.688	2p _y	-0.033	
H _{1y}	0.983	0.901	1.319	1s	0.107	57.8
H _{2y}	0.983	-0.901	1.319	1s	0.107	57.8
H ₃	1.884	0	0.057	1s	-0.002	-1.3
O	0	0	0	2p _y	0.712	

^a To satisfy the C_{2v} point group, each of the two methyl group conformations is chosen so that one hydrogen (H_3) is in the nodal plane of the oxygen 2p_y orbital, as shown in Figure 2b. ^b The average ^1H coupling is 38.1 G.

broadening can be attributed mainly to g anisotropy since the line shapes for each pair of $\pm M_I$ components are not inverted with respect to each other. Together with the presence of characteristic line shapes, this absence of an inversion center for the spectrum supports a spectral analysis in terms of different g -tensor components, the anisotropy from the ^1H hyperfine coupling contributing much less to the overall line width.

Analysis of the line shape structure in the first-derivative spectrum reveals that for each line of the septet, there is a clearly defined feature on the low-field side of the main component. This set of features is marked by asterisks in the second-derivative spectrum (Figure 1) and can be assigned to the seven hyperfine components associated with the maximum g -tensor component (g_1). Similarly, the high-field extrema of the main components, identified by the arrows marking the negative excursions in the first-derivative spectrum, are assigned to the pattern originating from the minimum g -tensor component (g_3). The spectral analysis is completed by assigning the major septet peaks in the second-derivative spectrum to the set of hyperfine lines belonging to the intermediate g -tensor component (g_2). As the sample temperature was raised above 100 K, the spectral features associated with the g_1 component gradually diminished in intensity until they had almost disappeared at 155 K (see supplementary material). The asymmetric line shape of the main absorption envelope remained, however, indicating that the motion of the radical at this temperature suffices only to bring about a partial averaging of the g -tensor components. The ESR parameters derived from the analysis of the low-temperature spectrum (Figure 1) are given in Table I.

The assignment of the spectrum to the oxygen-centered π radical cation of dimethyl ether is unambiguous. First, the interaction with six equivalent ^1H nuclei is expected for the two methyl groups undergoing free rotation, and the relatively small ^1H hyperfine anisotropy (Table I) is characteristic of β -hydrogen couplings.¹³ Second, the fact that the isotropic coupling, $a_\beta(6\text{H}) = 43.0$ G, greatly exceeds the corresponding value of 24.7 G for the six β hydrogens in the isoelectronic $(\text{CH}_3)_2\text{CH}$ radical¹⁴ is consistent with recent studies on alkoxy radicals¹⁵⁻¹⁷ which show that β -hydrogen couplings are significantly larger in oxygen-centered than in carbon-centered π radicals.¹⁸ Third, the appreciable positive g shift (Table I) for one of the g -tensor components (g_1) is typical of oxygen-centered radicals.¹⁵⁻¹⁷ Lastly, INDO calculations¹⁹ (Table II) confirm the structure of the radical (Figure 2) and yield an average value for $a_\beta(6\text{H})$ of 38.1 G, in reasonable agreement with the experimental result (43.0 G).

As depicted in Figure 2, the unpaired electron resides mainly in the p_y orbital on oxygen, the only other significant spin pop-

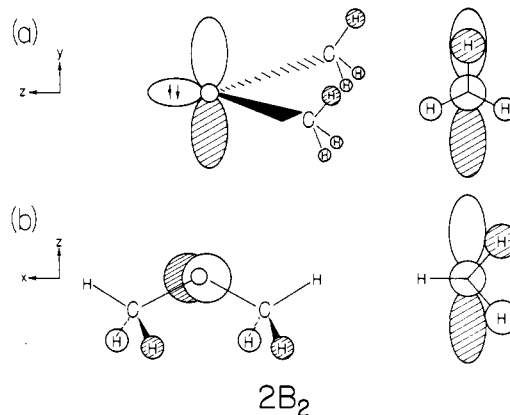


Figure 2. (a) Eclipsed and (b) bisected conformations of the methyl groups in the dimethyl ether radical cation, the CH_3 group orbitals possessing pseudo- π and π symmetry, respectively, with respect to the COC plane of the molecule. The oxygen p_y and the hydrogen π_y orbitals containing most of the spin population (Table II) are shown. Note that only conformation b possesses C_{2v} symmetry and that the calculated hydrogen 1s spin population for a rotating methyl group is obtained by averaging the results for the three hydrogens (Table II) in structure b.

ulations being in the hydrogen 1s orbitals (Table II) which constitute the methyl group orbitals. This $2B_2$ molecular orbital in the C_{2v} point group (Figure 2b) is usually described as a non-bonding orbital²⁰ since there is very little contribution from the p_y basis orbitals on the adjacent carbons. This situation arises because the p_y orbital on oxygen interacts with both the $\pi_y(\text{CH}_3)$ and $\pi_y^*(\text{CH}_3)$ group orbitals, resulting in the effective cancellation of the contributions from the methyl carbons.²¹ The g -tensor components can also be interpreted in terms of this structure for the radical. Thus, the large positive shift represented by g_1 is identified with the g_{xx} component since the magnetic field along this direction brings about the mixing of the singly occupied p_y orbital on oxygen with the $6A_1$ lone-pair orbital pointing along the z axis (Figure 2a), the energy separation between these orbitals being relatively small.^{20,22} A similar argument can be used to assign the intermediate g_2 component to g_{zz} , the p_y SOMO mixing in this case with the filled $4B_1(\sigma_{\text{CO}})$ MO immediately beneath the $6A_1$ level.²⁰ Finally, just as for the alkoxy radicals,¹⁶ the minimum g_3 (g_{yy}) component lies along the direction which is parallel to the symmetry axis of the singly occupied p orbital on oxygen.

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(18) The values of $a_\beta(6\text{H})$ obtained for the isoelectronic nitrogen-centered $(\text{CH}_3)_2\text{N}$ radical are 27.36 G³ and 28.48 G⁴. These results for β -hydrogen couplings from rotating methyl groups are much closer to those obtained for carbon-centered than for oxygen-centered radicals.¹⁵⁻¹⁷

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(21) A similar explanation is given in ref 20 (p 31) for the construction of the $2A''$ orbital in the eclipsed ethyl cation.

(22) The g_{max} component for alkoxy radicals is generally in the range 2.054-2.093¹⁶ and therefore is much larger than the value of 2.0138 obtained here for $(\text{CH}_3)_2\text{O}^+$. This difference is attributable to the likelihood of a much smaller energy separation between the two nonbonding oxygen orbitals in alkoxy radicals. These oxygen p orbitals are, of course, degenerate in isolated RO radicals and the degeneracy is lifted only by the asymmetric influence of hydrogen bonding in the solid state, a far smaller perturbation on the orbital energies than that produced in going from RO to R_2O^+ .

